

Lecture 3

Last time:

- categories of framed tangles
framed diagrams
- framed diagrams and framed tangles colored by a ribbon category $\mathcal{C}$
- the evaluation functor.

Quantum $\mathfrak{sl}_2$

For the double construction see notes from the last semester. Quantum $\mathfrak{sl}_2$ is a quotient algebra of the Drinfeld double of its quantum Borel subalgebra.

• $U_h \mathfrak{b}_+$ over $\mathbb{C}[[h]]$:

- generated by H and X with defining relations

$$[H, E] = 2E$$

complete over $\mathbb{C}[[h]]$, i.e.

as a vector space is the space of formal power series $\sum_{n=0}^{\infty} c_n(H, E) h^n$

where $C_n(H, E)$ is a polynomial in H & E .
 $n=0$

- it is a Hopf algebra with the comultiplication Δ acting on generators as

$$\Delta H = H \otimes 1 + 1 \otimes H, \quad \Delta E = E \otimes e^{\frac{1}{2}H} + 1 \otimes E$$

• Its restricted dual Hopf algebra $(U_h b_+)^{\vee}$ is generated by $H^{\vee}, E^{\vee}$ with defining relations

$$[H^{\vee}, E^{\vee}] = -\frac{1}{2} E^{\vee}$$

and with the comultiplication

$$\Delta H^{\vee} = H^{\vee} \otimes 1 + 1 \otimes H^{\vee}, \quad \Delta E^{\vee} = E^{\vee} \otimes e^{-2H^{\vee}} + 1 \otimes E^{\vee}$$

as a vector space it is $\text{Pol}(x)[H^{\vee}, h]$

The pairing with $U_h b_+$:

$$\langle (H^{\vee})^n (E^{\vee})^m, H^{n'}, E^{m'} \rangle = \delta_{nn'} \delta_{mm'}$$

$$n! (m)!,$$

$$(m)! = \frac{1 - e^{-hm}}{1 - e^{-h}} \cdots \frac{1 - e^{-h}}{1 - e^{-h}},$$

- The double of $U_h b_+$, based on the dual pair of Hopf algebras $U_h b_+$, $U_h b_+^\vee$ is the Hopf algebra generated by $H, X, H^\vee, X^\vee$. As a vector space, it is the space of polynomials in $H, X, X^\vee$ complete over formal power series in $H^\vee, h$ with defining relations

$$[H, H^\vee] = 0, [H, E] = 2E, [H^\vee, E] = \frac{h}{2}E,$$

$$[H, E^\vee] = -2E^\vee, [H^\vee, E^\vee] = -\frac{h}{2}E^\vee$$

$$[E, E^\vee] = e^{\frac{hH}{2}} - e^{-2H^\vee}$$

The comultiplication acts on generators as

$$\Delta H = H \otimes 1 + 1 \otimes H, \quad \Delta H^\vee = H^\vee \otimes 1 + 1 \otimes H^\vee,$$

$$\Delta E = E \otimes e^{\frac{hH}{2}} + 1 \otimes E, \quad \Delta E^\vee = E^\vee \otimes 1 + e^{-2H^\vee} \otimes E^\vee$$

This Hopf algebra is quasitriangular with the universal R -matrix

$$R = \sum_{n, m \geq 0} \frac{1}{n! (m)!} H^n E^m \otimes (H^\vee)^n (E^\vee)^m$$

$$= \exp(H \otimes H^\vee) f(E \otimes E^\vee; e^{-h}) \in \mathcal{D}(U_h \mathfrak{gl}_2)^{\hat{\otimes}}$$

$$f(z, q) = \sum_{m \geq 0} \frac{(1-q)^m z^m}{(1-q) \cdots (1-q^m)} = \prod_{n \geq 0} (1 - (1-q)zq^n)^{-1}$$

Here $\hat{\otimes}$ is the tensor product completed over $\mathbb{C}[[h, H^\vee \otimes 1, 1 \otimes H^\vee]]$.

- Define the algebra $U_h \mathfrak{gl}_2$ as the division algebra for $\mathcal{D}(U_h \mathfrak{gl}_2)$ where $H^\vee$ and $E^\vee$ are divisible by h . Define

$$G = \frac{4H^\vee}{h}, \quad F = \frac{E^\vee}{e^{h/2} - e^{-h/2}}$$

i.e. $U_h \mathfrak{gl}_2$ is generated by H, G, E, F over $\mathbb{C}[[h]]$, complete over $\mathbb{C}[[h]]$ with

$$[H, G] = 0, \quad [H, E] = 2E, \quad [H, F] = -2F$$

$$[G, E] = 2E, \quad [G, F] = -2F, \quad [E, F] = \frac{e^{\frac{hH}{2}} - e^{-\frac{hH}{2}}}{e^{h/2} - e^{-h/2}},$$

with

$$\Delta H = H \otimes 1 + 1 \otimes H, \quad \Delta G = G \otimes 1 + 1 \otimes G,$$

$$\Delta E = E \otimes e^{\frac{\hbar H}{2}} + 1 \otimes E, \quad \Delta F = F \otimes 1 + e^{-\frac{\hbar H}{2}} \otimes F,$$

and with

$$R = \exp\left(\frac{\hbar}{4} H \otimes G\right) \sum_{n \geq 0} \frac{(e^{\frac{\hbar}{2}} - e^{-\frac{\hbar}{2}})^n}{(n)!} E^n \otimes F^n$$

- Finally, it is clear that $H-G$ generates a Hopf ideal (just as it generates a Lie ideal in $\mathfrak{gl}_2$). The quotient algebra over this ideal is $U_{\hbar} \mathfrak{sl}_2$. it is generated by H, E, F with defining relations

$$[H, E] = 2E, \quad [H, F] = -2F$$

$$[E, F] = \frac{\text{sh}\left(\frac{\hbar}{2} H\right)}{\text{sh}\left(\frac{\hbar}{2}\right)}$$

with the comultiplication

$$\Delta H = H \otimes 1 + 1 \otimes H, \quad \Delta E = E \otimes e^{\frac{\hbar H}{2}} + 1 \otimes E, \quad \Delta F = F \otimes 1 + e^{-\frac{\hbar H}{2}} \otimes F$$

with the antipode

$$S(H) = -H, \quad S(E) = -E e^{-\frac{\hbar H}{2}}, \quad S(F) = -e^{\frac{\hbar H}{2}} F$$

and it is quasitriangular with

$$R = \exp\left(\frac{\hbar}{4} H \otimes H\right) \sum_{n \geq 0} \frac{(e^{\frac{\hbar}{2}} - e^{-\frac{\hbar}{2}})^n}{(n)!} E^n \otimes F^n$$

Important property:

Thm. 1) $U_{\hbar} \mathfrak{sl}_2 \simeq U \mathfrak{sl}_2[[\hbar]]$ as an algebra
(not as a Hopf algebra)

2) $Z(U_{\hbar} \mathfrak{sl}_2)$ is generated by

$$C_2 = \left(\frac{\text{sh}\left(\frac{\hbar}{4}(H+1)\right)}{\text{sh}\left(\frac{\hbar}{2}\right)} \right)^2 + FE$$

Finite dimensional representations of $U_{\hbar} \mathfrak{sl}_2$

let $U_{\hbar} \mathfrak{b}_+$ be the Borel subalgebra generated by H and E . Define L_{λ} , $\lambda \in \mathbb{C}$ as the one dimensional representation of $U_{\hbar} \mathfrak{b}_+$ with

$$H L_{\lambda} = 2\lambda L_{\lambda}, \quad X L_{\lambda} = 0$$

Define the Verma module M_λ over $U_h \mathfrak{sl}_2$ as the induced representation

$$M_\lambda = \text{Ind}_{U_h \mathfrak{b}_+}^{U_h \mathfrak{sl}_2} \mathbb{C}_\lambda$$

As a vector space

$$M_\lambda = U_h \mathfrak{sl}_2 \otimes_{U_h \mathfrak{b}_+} \mathbb{C}_\lambda \simeq U_h \mathfrak{n}_+ \otimes_{\mathbb{C}[\hbar]} \mathbb{C}_\lambda$$

where $U_h \mathfrak{n}_+$ is the subalgebra generated by F . Choose a basis e_λ^λ in $\mathbb{C}_\lambda$.

This gives a natural basis in M_λ

$$M_\lambda = \bigoplus_{n \geq 0} \mathbb{C} e_{\lambda - n}^\lambda$$

$$\text{where } e_{\lambda - n}^\lambda = F^n e_\lambda^\lambda.$$

Let us compute the action of generators on this basis. We immediately have

$$\gamma e_{\lambda-n}^{\lambda} = e_{\lambda-n-1}^{\lambda}$$

$$H e_{\lambda-n}^{\lambda} = H F^n e_{\lambda}^{\lambda} = F^n (H - 2n) e_{\lambda}^{\lambda} = 2(\lambda - n) e_{\lambda-n}^{\lambda}$$

$$E e_{\lambda-n}^{\lambda} = E F^n e_{\lambda}^{\lambda} = \left(\frac{\text{sh}(\frac{h}{2} H)}{\text{sh}(\frac{h}{2})} F^{n-1} + \dots \right.$$

$$\left. + F^k \frac{\text{sh}(\frac{h}{2} H)}{\text{sh}(\frac{h}{2})} F^{n-k-1} + \dots + F^{n-1} \frac{\text{sh}(\frac{h}{2} H)}{\text{sh}(\frac{h}{2})} \right) e_{\lambda}^{\lambda}$$

$$= \frac{1}{\text{sh}(\frac{h}{2})} \left(\text{sh}\left(\frac{h}{2}(2\lambda - 2n + 2)\right) + \text{sh}\left(\frac{h}{2}(\lambda - 2n + 4)\right) + \dots + \right. \\ \left. + \text{sh}\left(\frac{h}{2}\lambda\right) \right) e_{\lambda-n+1}^{\lambda}, \quad n \geq 1$$

i.e.

$$E e_{\lambda-n}^{\lambda} = \frac{\text{sh}\left(\frac{h}{2}(\lambda - n + 1)\right) \text{sh}\left(\frac{h}{2}n\right)}{\text{sh}\left(\frac{h}{2}\right)^2} e_{\lambda-n+1}^{\lambda}, \quad n \geq 1$$

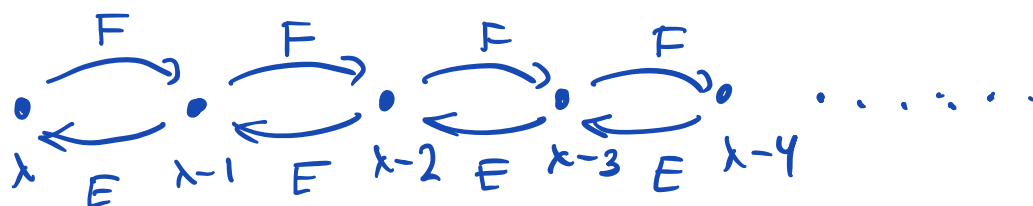
and, by definition of e_{λ}^{λ} , $E e_{\lambda}^{\lambda} = 0$.

The following is clear

Thm. M_λ is irreducible for generic complex λ (when $\lambda \neq 0, \frac{1}{2}, 1, \frac{3}{2}, 2, \dots$)

Indeed, it is easy to show that M_λ does not subrepresentations, except itself.

The action of H, E, F can be described by the diagram



Here vertices denote weight subspaces on which H acts diagonally.

However, when $\lambda = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$ there is an extra singular vector for E :

$$E e_{-\lambda-1}^\lambda = 0$$

As a consequence vectors $e_{-\lambda-1}^\lambda, e_{-\lambda-2}^\lambda, \dots$ form a subrepresentation $M_{-\lambda-1} \subset M_\lambda$.

The quotient representation

$$V_\lambda = M_\lambda / M_{-\lambda-1}$$

with the weight basis e_m^λ , $m = \lambda, \lambda-1, \lambda-2, \dots, -\lambda$ is the $2\lambda+1$ -dimensional irreducible $U_h \mathfrak{sl}_2$ -module.

Similarly to $\mathfrak{sl}_2$ these are all irreducible finite dimensional representations: any other is isomorphic to one of the V_λ 's.

Also similarly to $\mathfrak{sl}_2$ it is easy to show that any finite dimensional representation is isomorphic to a direct sum of irreducibles, i.e. for any finite dimensional $U_h \mathfrak{sl}_2$ -module

V we have

$$V \simeq \bigoplus_e V_e^{\oplus m_e(V)}$$

where $m_e(V)$ is the multiplicity of V_e in V .

Finally, it is easy to compute the action of the Casimir on the highest weight and therefore on the whole M_λ

$$c_2 M_\lambda = \left[\lambda + \frac{1}{2}\right]^2 M_\lambda$$

and therefore

$$c_2 V_e = \left[e + \frac{1}{2}\right]^2 V_e$$

Here

$$[e] = \frac{\text{sh}\left(\frac{h\rho}{2}\right)}{\text{sh}\left(\frac{h}{2}\right)}$$

The R-matrix in $V_{e_1} \otimes V_{e_2}$

Now we can compute the action of R on $V_{e_1} \otimes V_{e_2}$. Recall the action of generators of generators of $U_h \mathfrak{sl}_2$ on a weight basis in V_e :

$$H e_m = 2m e_m, \quad F e_m = e_{m-1},$$

$$E e_m = [e+m+1][e-m] e_{m+1},$$

For this action

$$E^n e_m = [e+m+1][e+m+2] \dots [e+m+n] \cdot \\ \cdot [e-m][e-m-1] \dots [e-m-n+1] e_{m+n}$$

$$F^n e_m = e_{m-n}$$

Let $(R)_{m_1, m_2}^{n_1, n_2}$ be matrix elements of R in the monomial basis $e_{m_1} \otimes e_{m_2}$:

$$R e_{m_1} \otimes e_{m_2} = \sum_{n_1, n_2} (R)_{m_1, m_2}^{n_1, n_2} e_{n_1} \otimes e_{n_2}$$

From the explicit formula for R we have:

$$R e_{m_1} \otimes e_{m_2} = \sum_{n \geq 0} \frac{e^{h \frac{n(n-1)}{4}}}{[n]!} \frac{[e+m_1+n]! [e-m_1]!}{[e+m_1]! [e-m_1-n]!} (e^{\frac{h}{2}} - e^{-\frac{h}{2}})^n \\ e^{h(m_1+h)(m_2-n)} e_{m_1+n} \otimes e_{m_2-n}$$

here we used the identity

$$(n) = \frac{1 - e^{-\hbar n}}{1 - e^{-\hbar}} = e^{-\frac{\hbar(n-1)}{2}} [n], \quad (n)! = e^{-\frac{\hbar}{2}(1+\dots+(n-1))} [n]!$$

$$= e^{-\frac{\hbar}{4}n(n-1)} [n]!$$

In particular when $\ell = \frac{1}{2}$,

$$R e_{\frac{1}{2}} \otimes e_{\frac{1}{2}} = e^{\frac{\hbar}{4}} e_{\frac{1}{2}} \otimes e_{\frac{1}{2}}$$

$$R e_{\frac{1}{2}} \otimes e_{-\frac{1}{2}} = e^{-\frac{\hbar}{4}} e_{\frac{1}{2}} \otimes e_{-\frac{1}{2}}$$

$$R e_{-\frac{1}{2}} \otimes e_{\frac{1}{2}} = e^{-\frac{\hbar}{4}} e_{-\frac{1}{2}} \otimes e_{\frac{1}{2}} + (e^{\frac{\hbar}{2}} - e^{-\frac{\hbar}{2}}) e^{-\frac{\hbar}{4}} e_{\frac{1}{2}} \otimes e_{\frac{1}{2}}$$

$$R e_{-\frac{1}{2}} \otimes e_{-\frac{1}{2}} = e^{+\frac{\hbar}{4}} e_{-\frac{1}{2}} \otimes e_{-\frac{1}{2}}$$

in the basis $\rightarrow$

we have

$$R = q^{\frac{1}{2}} \left(\begin{array}{c|c} q & \\ \hline 1 & \\ \hline q^{-1} & 1 \\ \hline & q \end{array} \right), \quad \begin{pmatrix} e_{\frac{1}{2}} \otimes e_{\frac{1}{2}} \\ e_{\frac{1}{2}} \otimes e_{-\frac{1}{2}} \\ e_{-\frac{1}{2}} \otimes e_{\frac{1}{2}} \\ e_{-\frac{1}{2}} \otimes e_{-\frac{1}{2}} \end{pmatrix}$$

This R -matrix gives Jones polynomial.

The decomposition of the tensor product

Let V_{ℓ_1} and V_{ℓ_2} be two finite dimensional $U_{\hbar} \mathfrak{sl}_2$ -modules. Their tensor product is finite dimensional and therefore is isomorphic to a direct sum of V_{ℓ} 's. It turns out that

$$(*) \quad V_{\ell_1} \otimes V_{\ell_2} \cong V_{\ell_1+\ell_2} \oplus V_{\ell_1+\ell_2-1} \oplus \dots \oplus V_{|\ell_1-\ell_2|}$$

For example, the irreducible submodule with the highest weight $\ell_1+\ell_2$ is easy to construct:

- $e_{\ell_1}^{\ell_1} \otimes e_{\ell_2}^{\ell_2}$ is a highest weight vector

(of weight $\ell_1+\ell_2$) in this tensor product,

i.e. $H(e_{l_1}^{l_1} \otimes e_{l_2}^{l_2}) = (l_1 + l_2) e_{l_1}^{l_1} \otimes e_{l_2}^{l_2}$ and any other highest weight vector is its multiple.

- The action of F on this vector generate the subrepresentation $V_{l_1+l_2} \subset V_{l_1} \otimes V_{l_2}$. It is easy to check that $F^{l_1+l_2+1}(e_{l_1}^{l_1} \otimes e_{l_2}^{l_2}) = 0$.

- The weight subspace of weight l_1+l_2-2 in this tensor product is two dimensional with the monomial basis $e_{l_1-2}^{l_1} \otimes e_{l_2}^{l_2}$ and $e_{l_1}^{l_1} \otimes e_{l_2-2}^{l_2}$. Vector

$$F(e_{l_1}^{l_1} \otimes e_{l_2}^{l_2}) = e_{l_1-2}^{l_1} \otimes e_{l_2}^{l_2} + e^{-h l_1} e_{l_1}^{l_1} \otimes e_{l_2-2}^{l_2}$$

is in $V_{l_2+l_2}^{l_1+l_2} \subset V_{l_1} \otimes V_{l_2}$. We can choose the second linearly independent vector in this weight subspace as a

highest weight vector of weight $\ell_1 + \ell_2 - 2$.

The equation

$$E \left(a e_{\ell_1-2}^{\ell_1} \otimes e_{\ell_2}^{\ell_2} + b e_{\ell_1}^{\ell_1} \otimes e_{\ell_2-2}^{\ell_2} \right) = 0$$

Fixes a and b uniquely up to a common multiple.

HW, Find $a:b$

Continuing this process one obtains an explicit decomposition of the tensor product into irreducible components.

Theorem [KR] Embeddings $V^\ell \hookrightarrow V^{\ell_1} \otimes V^{\ell_2}$

in the decomposition (*) in terms of weight basis e_m^ℓ

$$H e_m^\ell = 2m e_m^\ell, \quad F e_m^\ell = e_{m-1}^\ell, \quad F e_{-\ell}^\ell = 0$$

$$E e_m^\ell = \frac{\text{sh}\left(\frac{\hbar}{2}(\ell+m+1)\right) \text{sh}\left(\frac{\hbar}{2}(\ell-m)\right)}{\text{sh}\left(\frac{\hbar}{2}\right)^2} e_{m+1}^\ell$$

are given by formulae:

$$e_m^l = \sum_{m_1+m_2=m} \begin{bmatrix} l_1 & l_2 & l \\ m_1 & m_2 & m \end{bmatrix} z_l e_{m_1}^{l_1} \otimes e_{m_2}^{l_2}$$

where

$$\begin{bmatrix} l_1 & l_2 & l \\ m_1 & m_2 & m \end{bmatrix} = z_l (-1)^{j_1-m_1} \frac{\rho(x)^{\frac{1}{2}}}{dn} {}_3\phi_2 \left(\begin{matrix} q^{-n}, abq^{x+1}, q^{-x} \\ aq, q^{-N} \end{matrix} ; q, q \right)$$

$$n = j - m, \quad a = q^{j_2-j+m_1}, \quad b = q^{j_1-j_2+m}, \quad x = j_1 - m_1$$

$$N = j + j_2 - m_1, \quad q = e^h$$

$\rho(x)$, dn see [KR].

$${}_{p+1}\phi_p \left(\begin{matrix} a_1, \dots, a_{p+1} \\ b_1, \dots, b_p \end{matrix} ; q, z \right) = \sum_{k=0}^{\infty} \frac{(a_1, \dots, a_{p+1}; q)_k}{(b_1, \dots, b_p; q)_k} z^k,$$

$$(a; q)_k = \prod_{i=0}^{k-1} (1 - aq^i),$$

$$(a_1, \dots, a_p; q)_k = \prod_{j=1}^p (a_j; q)_k,$$

and z_l are arbitrary nonzero scalars.

Remark. The decomposition (*) means that $\text{Hom}(V_\ell, V_{\ell_1} \otimes V_{\ell_2})$ are 1-dimensional for $\ell_1 + \ell_2 \geq \ell \geq |\ell_1 - \ell_2|$ and zero otherwise. Different choices of z_ℓ correspond to different choices of the basis in $\text{Hom}(V_\ell, V_{\ell_1} \otimes V_{\ell_2})$.